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APPLICATION OF NUMERICAL METHODS TO THE CALCULATION OF ELECTROSTATIC FIELDS IN AIRCRAFT FUEL TANKS

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0. Abstract	
The solution of electrostatic field problems occur	rring during the
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and the finite element method is used to calculat	
potential distribution within the tank.	
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Deputy Commander

1. ELECTROSTATIC HAZARDS IN AIRCRAFT FUEL TANKS DURING REFUELLING

The generation of high levels of electrostatic charge in hydrocarbon fuels (such as JP-4) during aircraft refuelling has long been recognized as an explosion hazard. The fuel is pumped at high flow rates through pipes, hoses, and filter/separators, and hence is exposed to relatively large liquid/solid interfaces. The double layer created at these interfaces coupled with the movement of fuel across them leads to a net unipolar charge being acquired by the fuel as it is swept along. The high charging tendency of a fuel such as JP-4, coupled with a low conductivity (typically < 10 pS/m) can lead to hazardous charge accumulation in the receiving tank. If resulting local electrostatic fields on the fuel surface exceed the breakdown value for the vapour space, electrostatic discharges may occur. Such a discharge may be incendive if (i) it has sufficient energy, and (ii) the fuel/air mixture lies in a flammable range (i.e. will support combustion). It is thus patently of importance to be able to estimate electrostatic potential and field distributions in fuel tanks for given charge distributions.

The introduction of polyurethane foam into fuel tanks to act as an explosion suppressant presents an additional problem. The foam itself acts as a secondary charge generating surface for the fuel (Ref. 1). The relaxation time of the charge in the tank increases enormously, thus leading to substantially increased levels of accumulated charge in the tank, since the charge is unable to relax to earth.

The present study is an attempt to apply numerical and computational techniques to the problem in order to examine the feasibility of providing useful working guides for estimating electrostatic potentials and fields in fuel tanks containing such foams.

2. THE MATHEMATICAL PROBLEM AND THE FINITE ELEMENT METHOD

One method of approach to the problem is to consider the fuel tank at various levels of filling, postulate a charge distribution and boundary conditions, and solve Poisson's equation for this situation. This approach has been tackled analytically for some very simple geometries and charge distributions (Refs. 2,3). For the modelling of realistic situations, however, the analysis becomes intractable, and recourse to numerical methods becomes essential.

Various numerical techniques for approximating the solution of electrostatic field problems are currently in use. These include the finite-difference method, the charge simulation method, the boundary integral method, and the finite element method. The method chosen for this study is the finite element method. It allows the modelling of complicated geometries, inhomogeneous charge distributions, and dielectric changes within the region of interest. Its disadvantages are relatively large data input and large computer storage requirements.

The basic mathematical problem to be tackled is the solution of Poisson's equation

$$\nabla^2 \phi = -\frac{\rho}{\varepsilon_0 \varepsilon_r}$$

within a region, subject to certain boundary conditions. Here, ϕ is the electrostatic potential, ρ is the space charge density, ϵ_{α} is the

absolute permittivity (8.854 x 10^{-12} F/m) and ϵ_{r} is the relative dielectric constant of the medium. To provide a complete definition of problem, one or more of the following boundary conditions are required:

- (i) $\phi = f(s)$ which fixes the potential at the boundary s, as a specified function f(s) of position. (e.g. on an earthed boundary $\phi = 0$)
- (ii) $\frac{\partial \phi}{\partial n} = 0$ which forces equipotentials to cross a boundary normally
- (iii) $\frac{\partial \phi}{\partial n} = h(s)$ which superimposes a surface charge density distribution on a boundary.

Let us now consider the functional

$$\mathbf{F} = \int_{\mathbf{V}} \frac{1}{2} \left\{ \left| \nabla \phi \right|^2 - \frac{20}{\varepsilon_0 \varepsilon_r} \phi \right\} d\mathbf{V} + \int_{\mathbf{S}} h \phi d\mathbf{S}$$

where V is the volume of the region of interest contained by the bounding surface S. If we minimise this functional, i.e. look for admissible potential functions ϕ such that

$$\delta \mathbf{F} = 0$$

it may be shown that the ensuing Euler-Lagrange equation is

$$\nabla^2 \phi = -\frac{\rho}{\varepsilon_0 \varepsilon_r}$$

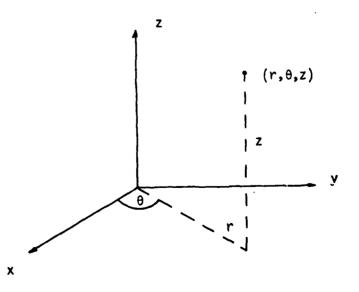
i.e. the potential function must satisfy Poisson's equation, subject to appropriate boundary conditions.

Briefly, the finite element method utilizes the above results by subdivision of the region of interest into a collection of elements, and approximating the potential by a set of piecewise continuous functions on these. Each element has a fixed number of nodes, and the minimisation is performed in each element with respect to the potential values at these nodes.

For a full three-dimensional problem, however, the resulting set of linear equations in the nodal potentials to be solved is generally very large, and necessitates substantial computing resources. For many problems, however, including the one under consideration, the problem size may be reduced in size and complexity by using special features of the geometry. It will be seen that for the fuel tank under consideration, the geometry may be taken to be axisymmetric, yielding results which are sufficiently accurate to make the approximation acceptable. To exploit axisymmetry, cylindrical coordinates are adopted, and Poisson's equation becomes, in a usual notation,

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial\phi}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2\phi}{\partial\theta^2} + \frac{\partial^2\phi}{\partial z^2} = -\frac{\rho}{\varepsilon_0\varepsilon_r}$$
(1)

where r, θ , z are cylindrical coordinates as shown:

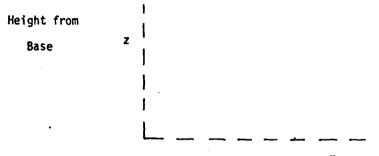


Page 4

For axisymmetric problems, the geometry and charge distributions are rotationally symmetric, the potential ϕ , hence rotationally symmetric, and therefore independent of θ . Poisson's equation thus reduces to

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \phi}{\partial r} \right) + \frac{\partial^2 \phi}{\partial z^2} = -\frac{\rho}{\varepsilon_0 \varepsilon_r}$$
 (2)

The problem is thus basically two-dimensional:



Radial distance from axis of symmetry

The finite element method consists of four basic steps. Firstly a grid of numbered nodal points is established over the region of interest, including the boundary. At each interior node the value of potential is to be determined. On the boundary ϕ or $\frac{\partial \phi}{\partial n}$, its normal derivative, is given. Secondly the nodes are interconnected to form a finite number of subregions which collectively approximate the region of interest.

Thirdly the potential is approximated by a continuous function over each subregion, continuity conditions being imposed at subregion boundaries. Finally the unknown values of the potential at the node points are calculated by means of a variational principle, e.g. minimisation of electrostatic energy.

The subregions used in this case are triangular. A first order interpolating polynomial (3) is used to approximate the potential distribution over each triangle.

$$\phi = \alpha_1 + \alpha_2 r + \alpha_3 z \tag{3}$$

Thus in a triangle whose nodes are labelled i, j, k we have -

$$\phi_{i} = \alpha_{1} + \alpha_{2}r_{i} + \alpha_{3}z_{i}$$

$$\phi_{j} = \alpha_{1} + \alpha_{2}r_{j} + \alpha_{3}z_{j}$$

$$\phi_{k} = \alpha_{1} + \alpha_{2}r_{k} + \alpha_{3}z_{k}$$

The variables α_1 α_2 α_3 can therefore be obtained in terms of ϕ_i ϕ_j ϕ_k and the coordinates of the noes i,j,l. Substituting back into the interpolating polynomial (3) gives -

$$\phi = N_{i}\phi_{i} + N_{j}\phi_{j} + N_{k}\phi_{k}$$
 (4)

where N_i, N_j, N_k are functions of r_i, r_j, r_k, z_i, z_j, z_k, r and z. $\frac{\partial \varphi}{\partial r} \text{ and } \frac{\partial \varphi}{\partial z} \text{ may also be readily obtained.}$

The unknown values of may now be obtained by minimising the functional -

$$F = \int_{A} \left\{ r \left(\frac{\partial \phi}{\partial r} \right)^{2} + r \left(\frac{\partial \phi}{\partial z} \right)^{2} - \frac{2\rho r \phi}{\varepsilon_{0} \varepsilon_{r}} \right\} . dA - 2 \int_{C} \phi h(s) . ds$$

i.e. $\delta F = 0$

It is not difficult to show that this minimisation is equivalent to the requirement that ϕ satisfies the axisymmetric Poisson equation with appropriate boundary conditions.

Substitution of the appropriate local expressions for ϕ , $\frac{\partial \phi}{\partial r}$ and $\frac{\partial \phi}{\partial z}$ into the functional, and minimisation by a Rayleigh-Ritz or equivalent technique yields a set of simultaneous equations in the nodal potentials for each triangle. These elemental equations are then assembled to give a global set of equations which may be solved using standard techniques.

3. DESCRIPTION OF COMPUTATIONAL MODEL

(i) Geometry

The computational model used is based on the drawings of an A-10 fuel tank (Ref. 4). The real tank configuration (Figure 1) is such as to permit an axisymmetric approximation. The axisymmetric configuration chosen is shown in Figure (2), cut away to show the inlet nozzle and the explosion suppressant foam blocks. A dimensional cross-section of the tank is shown in Figure (3). (Dimensions are in millimetres).

(ii) Boundary Conditions

Figure (4) illustrates the assumed boundary conditions. The tank walls and the inlet nozzle are assumed to be at earth potential $(\phi = 0)$.

Conforming to the real situation, the foam is divided into two categories, fixed and removable. Section 4 in Figure (4) corresponds to the fixed foam region. Sections 1,2, and 3 correspond to removable foam target sections. Thus, possible target configurations are Section 1 only, Sections 1 and 2 together, or Sections 1,2, and 3 together, allowing alteration of the voiding volume.

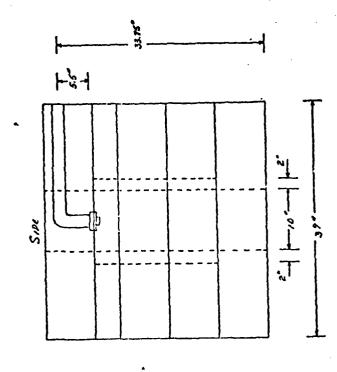
The dielectric boundaries of the foam blocks are also shown in Figure (4), as is the axis of symmetry. Since the problem is axisymmetric, the potential ϕ must satisfy $\frac{\partial \phi}{\partial n} = n$ on the axis of symmetry, i.e. the equipotentials must cross the axis normally.

(iii) The Charge Density

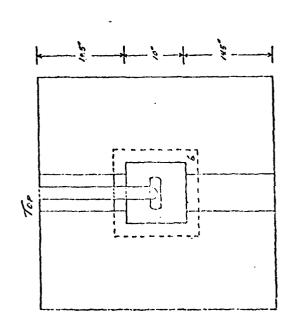
Accurate estimations of the volume and surface charge densities occurring when fuel is pumped at speed into a receiving tank are difficult

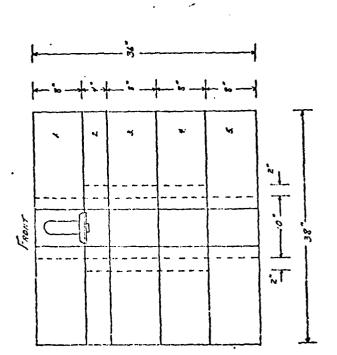
YOIDING CONCERT ILLUSTRATES FUTURE POLYETHER
BLUE COARSE PORE FORM CATTERIA
FOR THE F-4 AIRCARFT

B. MIDDLE PIECE OF FOAM WITH A LATINGS FOID AND A 6"8" 15 5" "REMANDLE A. BOTTOM PIECE OF FOAM WITH A 10" A 10" WID AND A 6" 8" A 8" A 8" REMOINSLE 6. BOTTOM FIECE OF FOAM WITH A 10" A 10" 8" VOID AND A 6" A 8" A 145" REMOINSLE 6. THREET FOAM SECTION - 14" A 14" A 50" WITH A 10" A 10" A 10" WID.



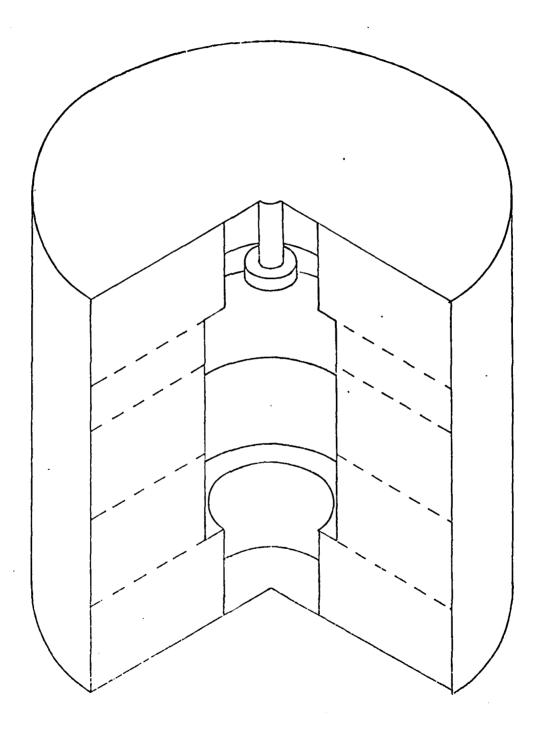
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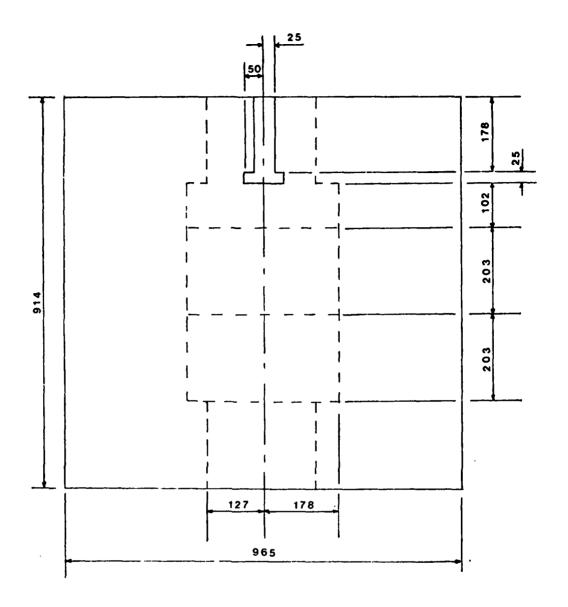
A-10 Fuel Tank

Fig. 1.



Isometric Projection of Tank Model

fig. 2.

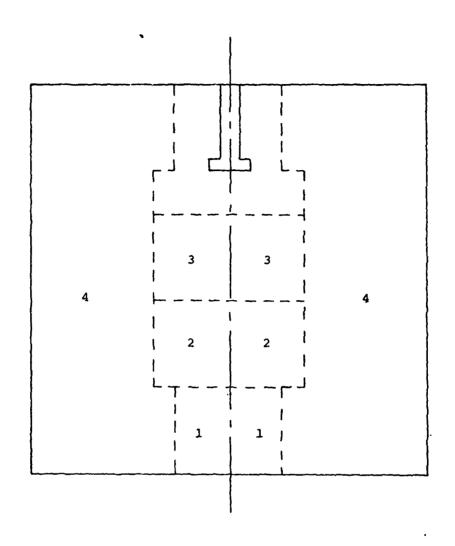


All Dimensions in mm.

Dimensioned Section of Tank Model

fig. 3.

Page 10



Earthed Boundaries

---- Dielectric boundaries

Axis of symmetry

1 Target section 1
2 Target section 2
3 Target section 3
4 Fixed foam region

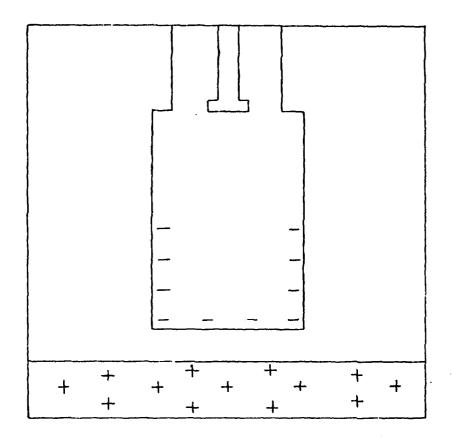
Section of Tank Model showing Target Sections

fig. 4.

to obtain. Factors such as fuel conductivity, rate of flow, turbulence, all contribute to the distribution of charge in the fuel at any particular time. The simplest, and not unrealistic, distribution to assume is a constant charge density within the fuel. There is evidence (Ref. 5) that jet fuels such as JP-4 generate static electricity when passing through porcus media. The explosion suppressant foam acts as a secondary static charge generating surface. Furthermore, the movement of charge with foam present is very slow, so a constant charge density at any filling level is a reasonable approximation. The dielectric constant of both fuel and foam has been taken to be 2. A typical charge distribution considered s shown in Figure (5). As we shall see, the charge distributions studied at various filling levels allow a wide choice of postulated charge distribution.

4. RESULTS

The cases studied consist of charge distributions, both volume and surface, for various filling levels, and for different target sections. A 'standard' constant volume charge distribution of 10⁻⁴ C/m³ in the fuel was chosen, together with a 'standard' surface charge distribution of 10⁻³ C/m² on the foam surface in the voiding region. This is not restricted since the two charge distributions are treated separately, and the principle of superposition allows simple scaling to a desired charge density. The plots provided give ten equipotentials between minimum and maximum potentials for each situation. Also given are plots of electrostatic field along the axis of symmetry for each configuration.



Typical Charge Distribution

Fig.5.

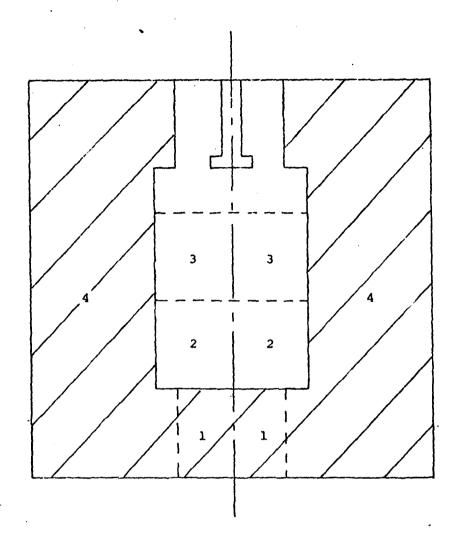
PLOTS 1 - 17

The configuration for plots 1-17 consists of foam Section 4 together with foam Section 1 inserted. (See Figure 6).

- Plots 1 5 Charge density in fuel $\approx 10^{-4}$ C/m³

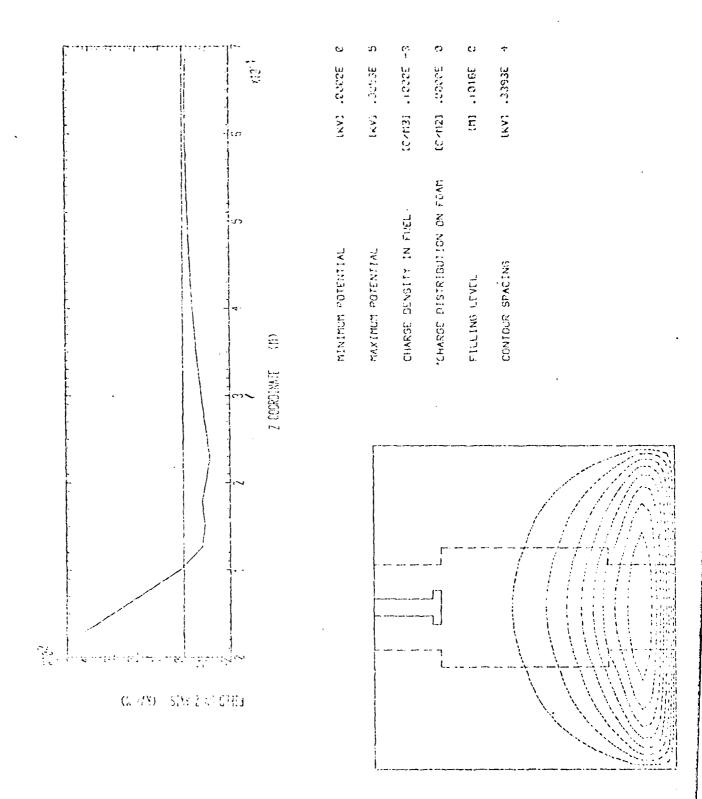
 Charge density on foam surface ≈ 0 C/m²

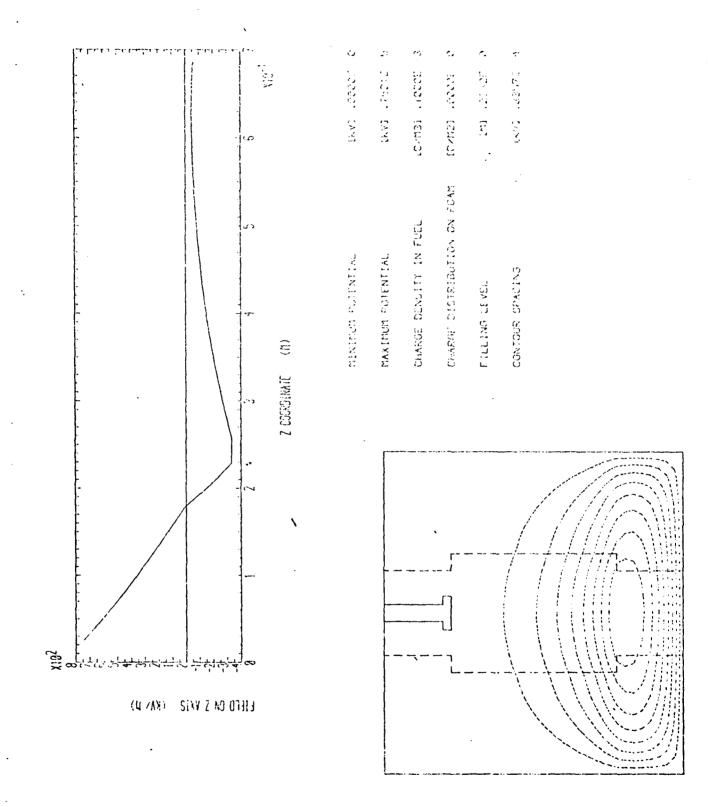
 Filling levels at .1, .2, .3, .4, .5 metres above base.
- Plot 6 Surface charge density of 10⁻³C/m² on Section 1 upper surface only.
- Plots 7 11 Surface charge on Section 1 upper surface and at heights .1, .2, .3, .4, .5 metres above this surface.
- Plot 12 Vertical stream of fuel with charge density 10⁻⁴ C/m³ impinging on the target area.
- Plots 13 17 Vertical stream of fuel + voiding region filled to heights .1, .2, .3, .4, .5 metres.



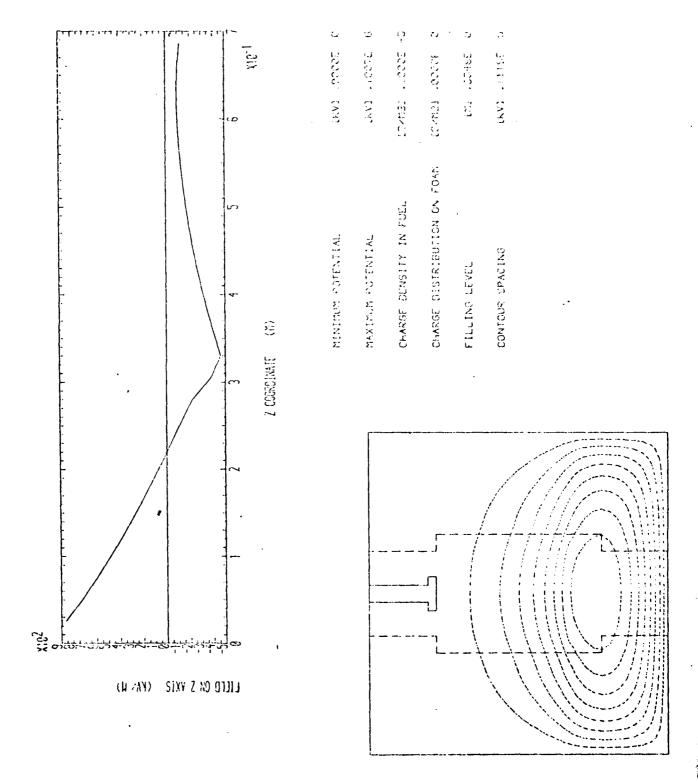
Foam sections 1 and 4 in place

fig. 6.

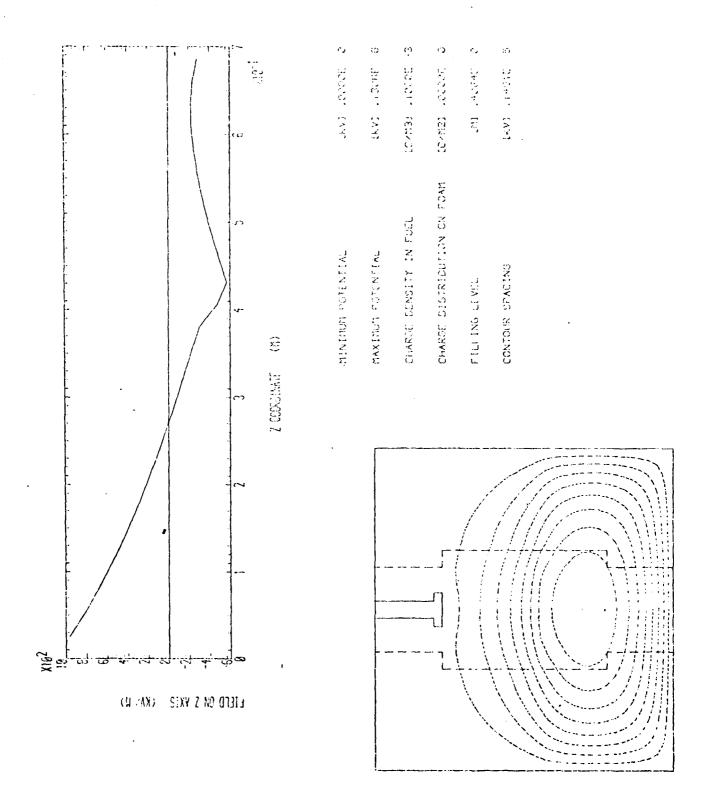




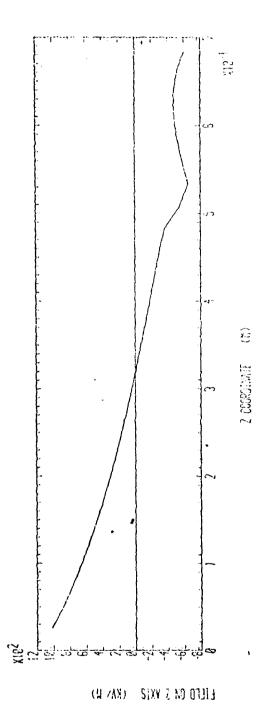
PLOT 2

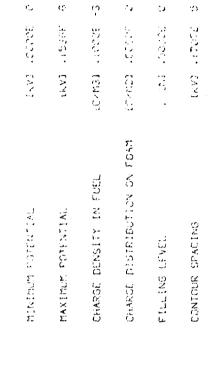


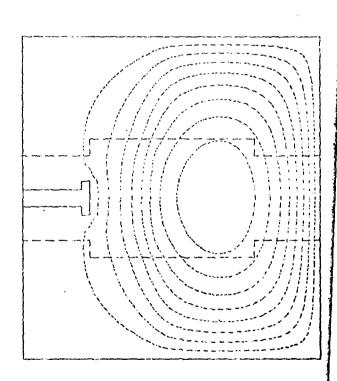
PLOT 3



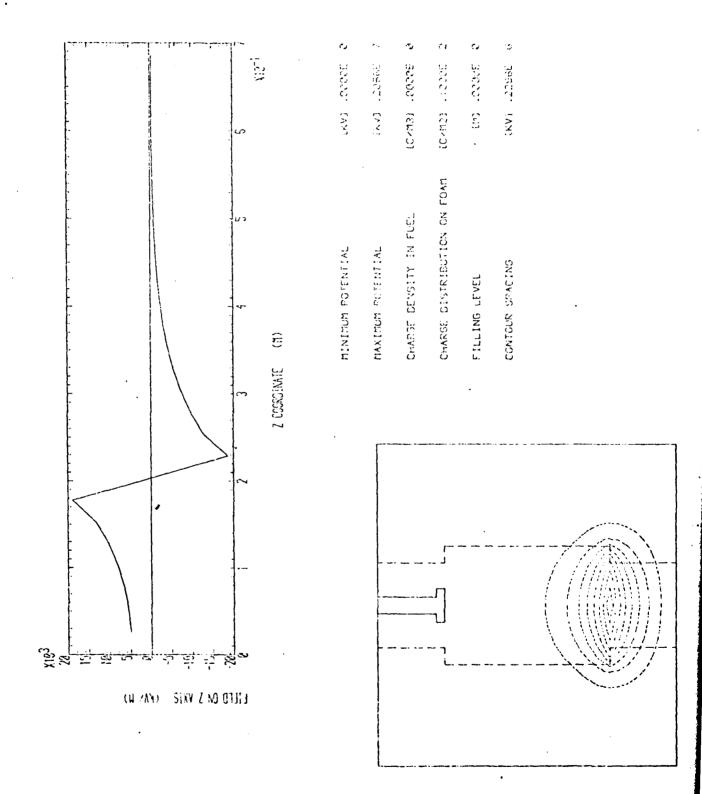
PLOT 4



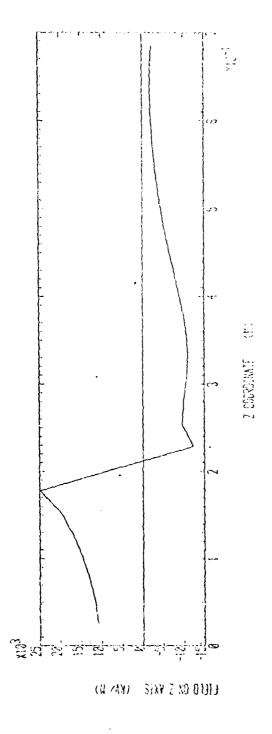




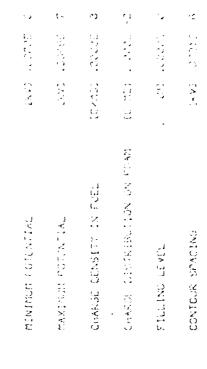
PLOT 5

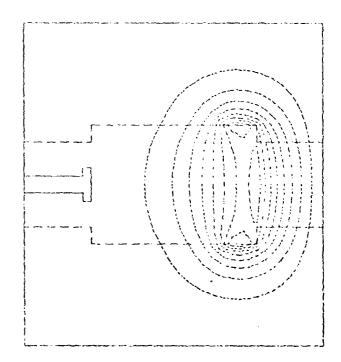


PLOT 6

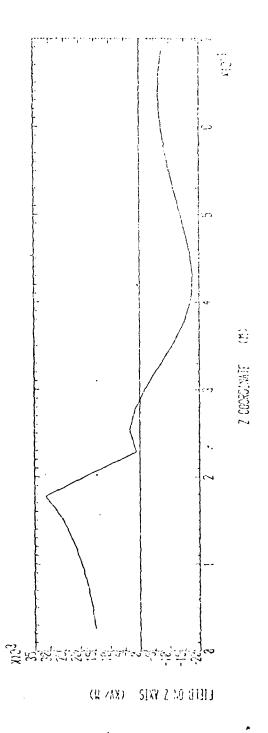


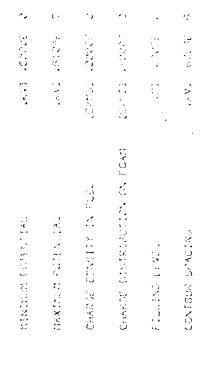
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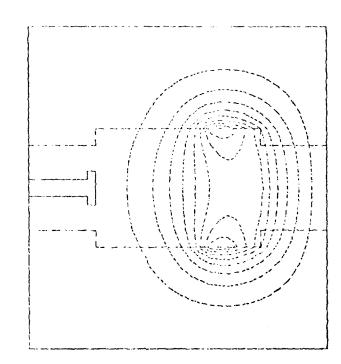




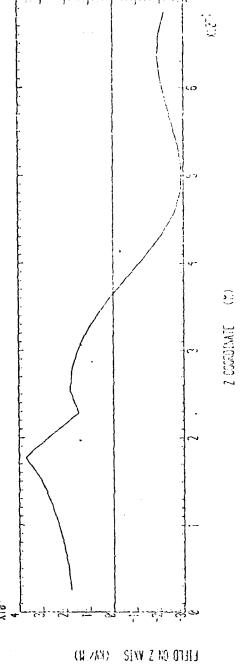
PLOT 7



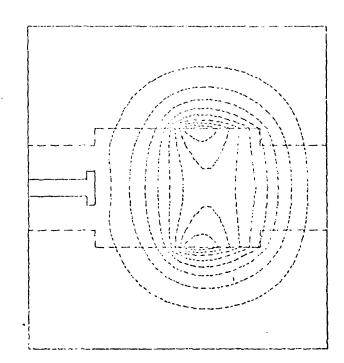




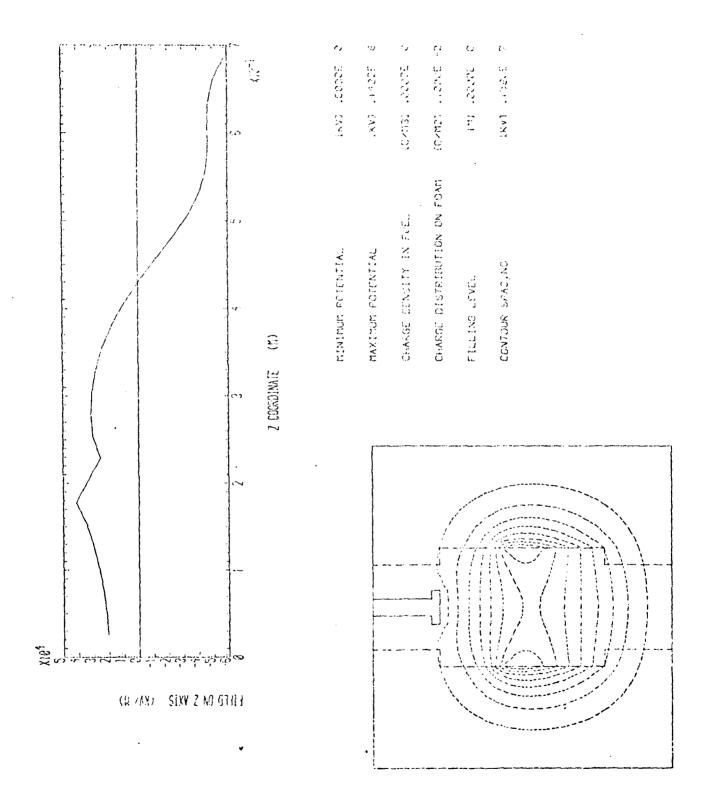
PLOT 8



MAKIMUM POTENTIAL CHAROE CENSITY IN FUEL CHAROE DESINIBUTION ON FOAM	Beret Card	TOWER TWO TENDER A SERVER A SE	
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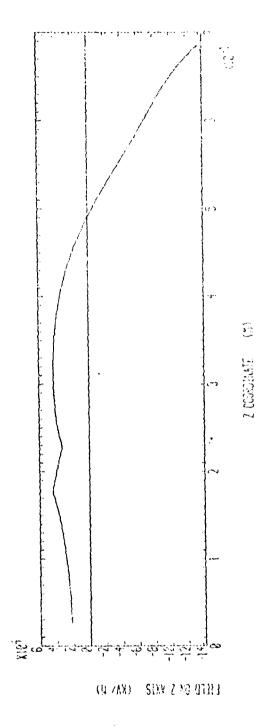


PLOT 9

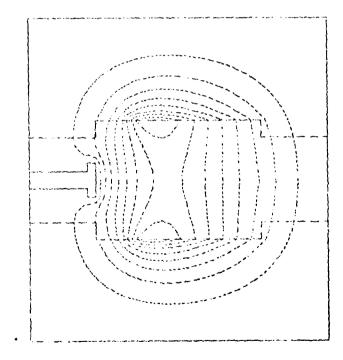


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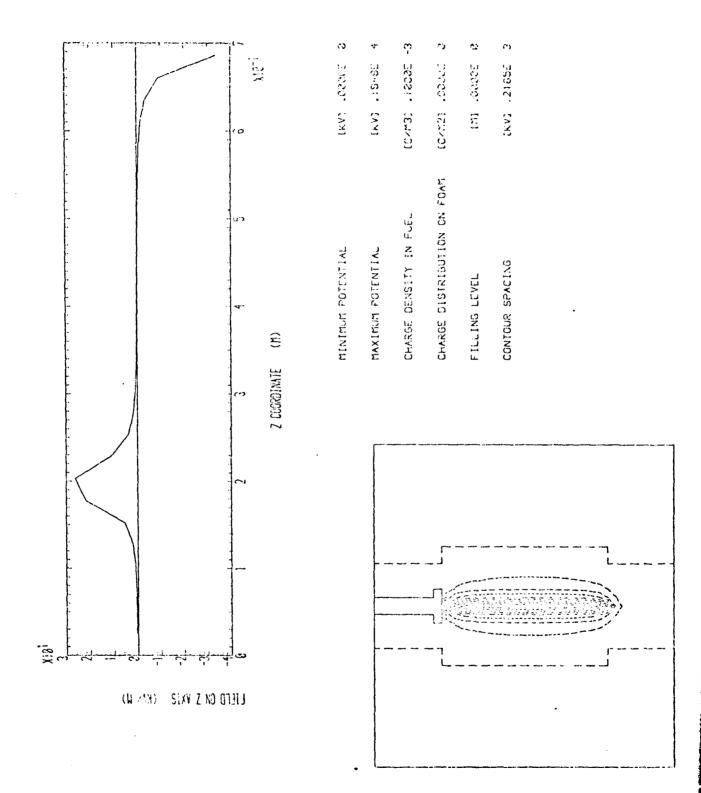
PLOT 10



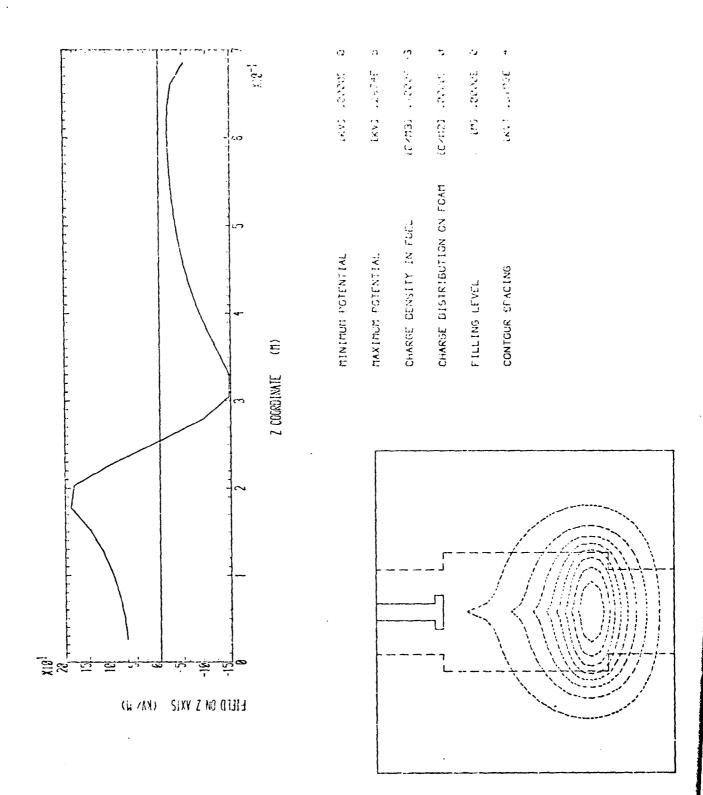




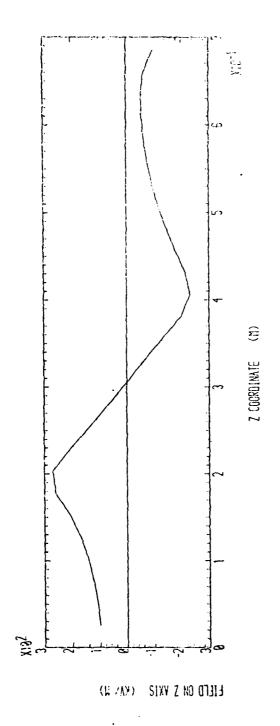
PLOT 11

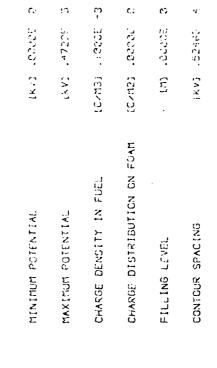


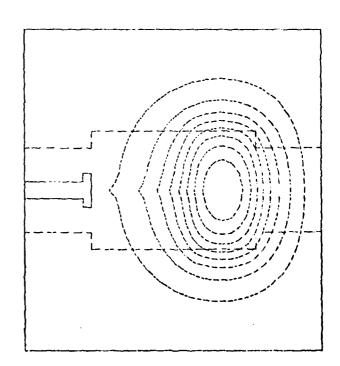
PLOT 12



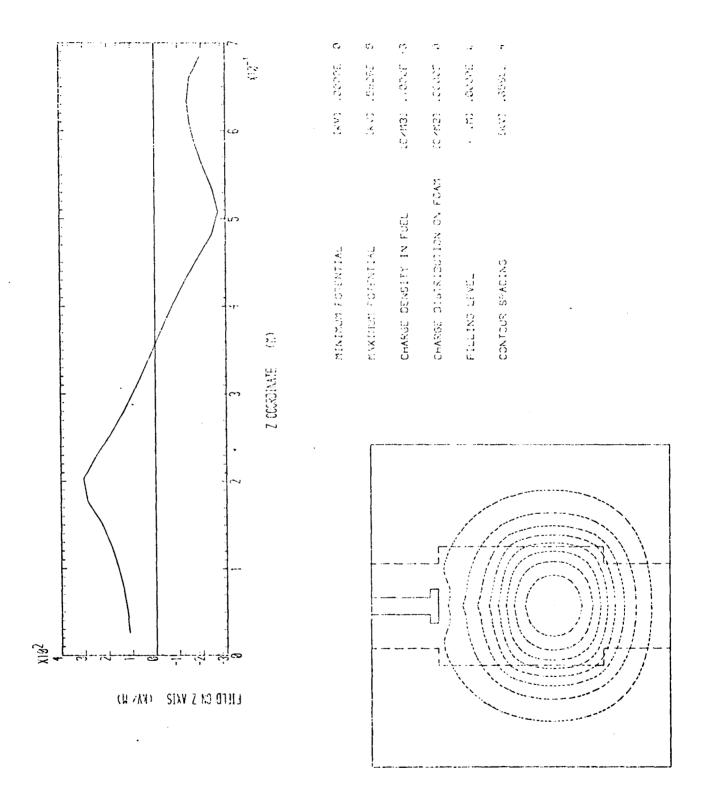
PLOT 13



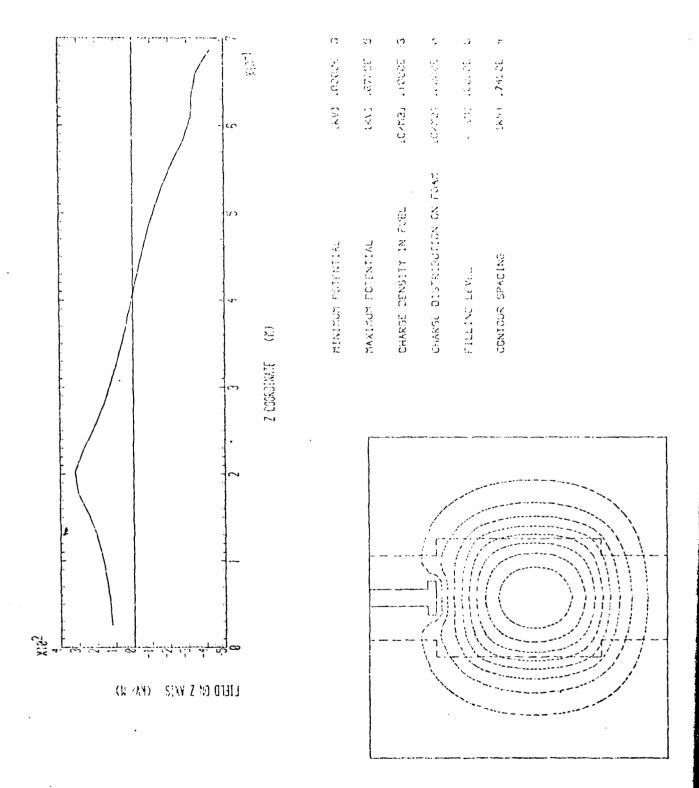




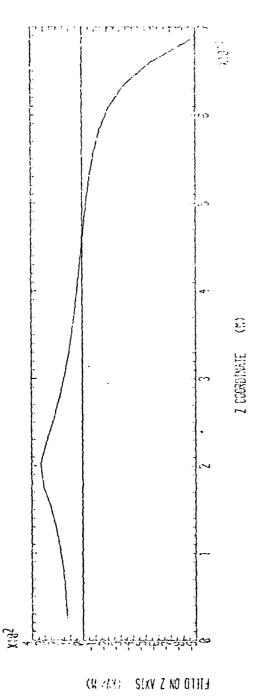
PLOT 14

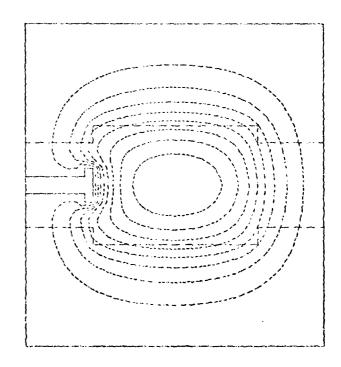


PLOT 15



PLOT 16





PLOT 17

PLOTS 18 - 30

The configuration for plots 18 - 30 consists of foam Section 4 together with foam Sections 1 and 2 inserted. (See Figure 7).

Plots 18 - 22	Charge	density	in	fuel			=	: 10)-4	C/m³
	Charge	density	on	foam	sui	face	=	•	0	C/m²
	Filling	levels	at	.1,	.2,	.3,	.4,	•5	met	res
	above b	ase.								

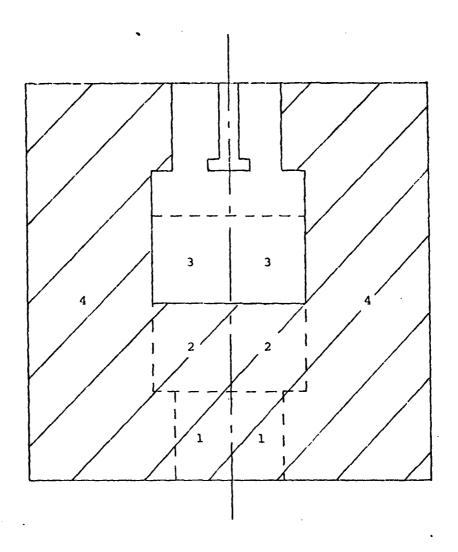
Plot 23 Surface charge density of 10⁻³ C/m² on Section 2 upper surface only.

Plots 24 - 26 Surface charge on Section 2 upper surface and at heights .1, .2, .3 metres above this surface.

Plot 27 Vertical stream of fuel with charge density

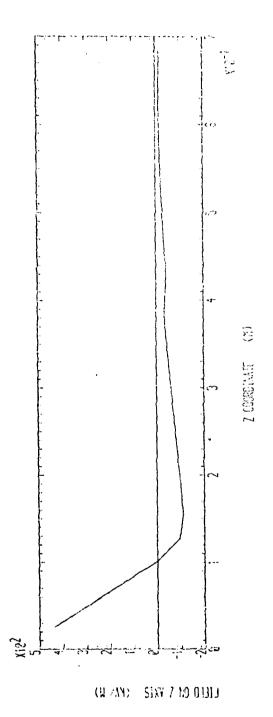
10⁻⁴ C/m³ impinging on the target area.

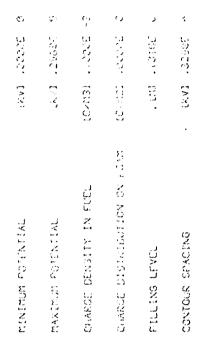
Plots 28 - 30 Vertical stream of fuel + voiding region filled to heights .1, .2, .3 metres.

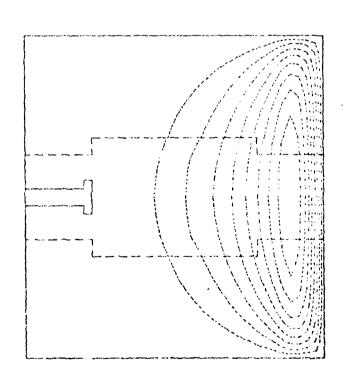


Foam sections 1 , 2 , and 4 in place

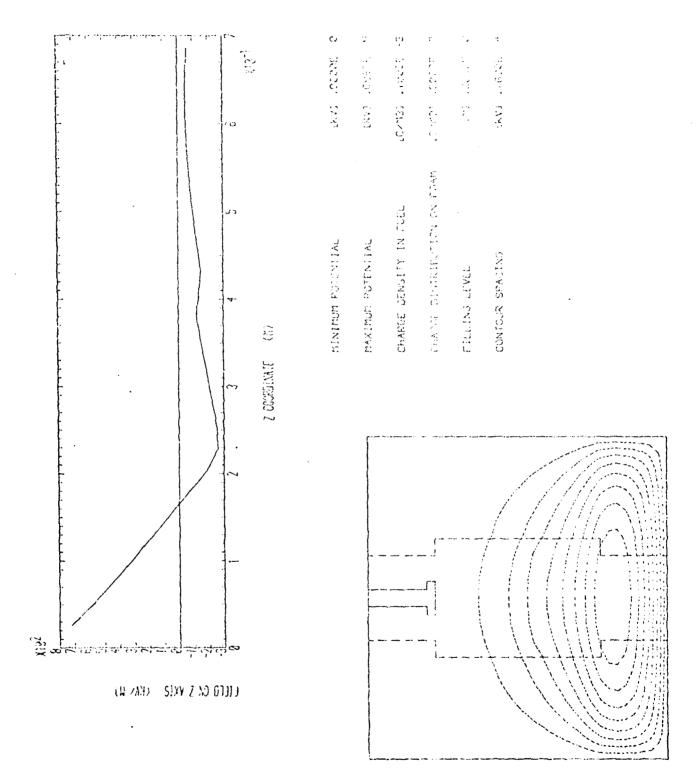
fig. 7.



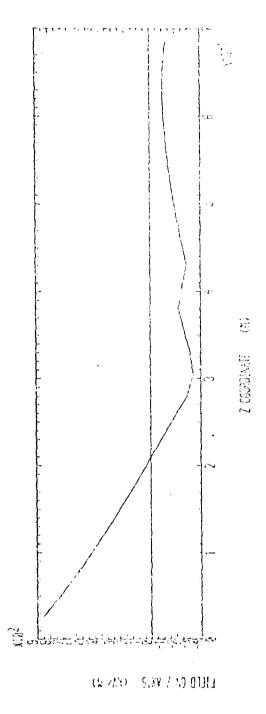


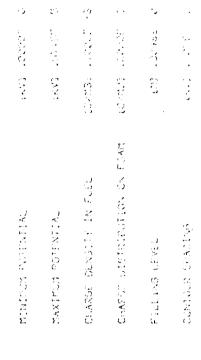


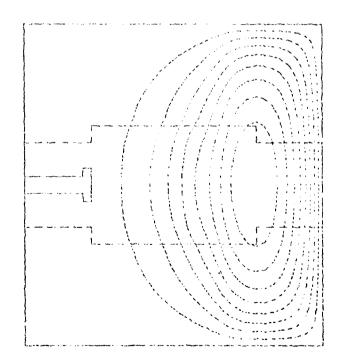
PLOT 18



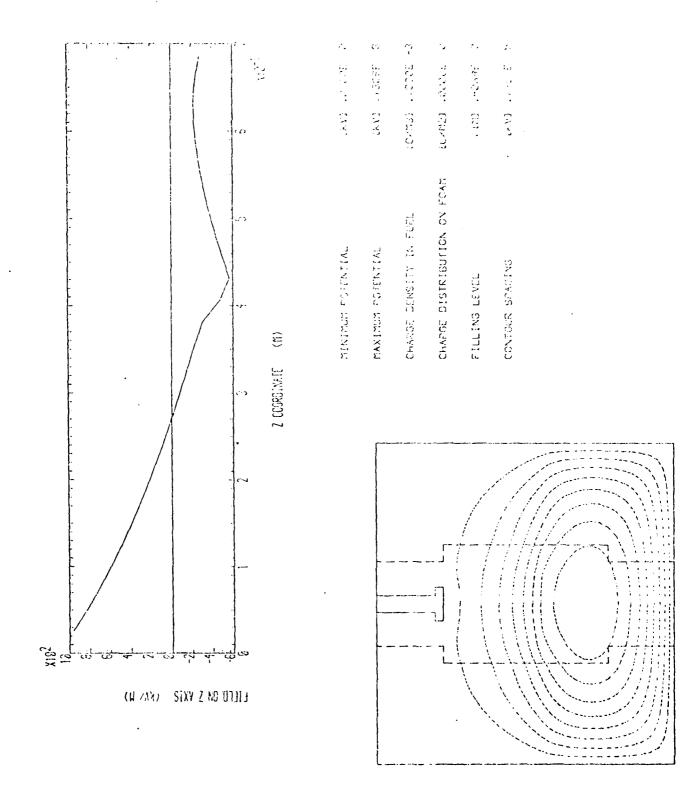
PLOT 19



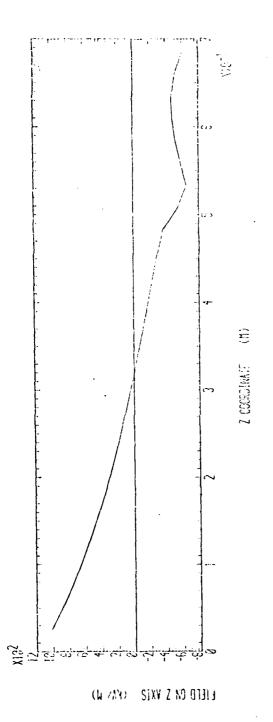


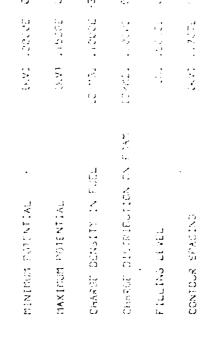


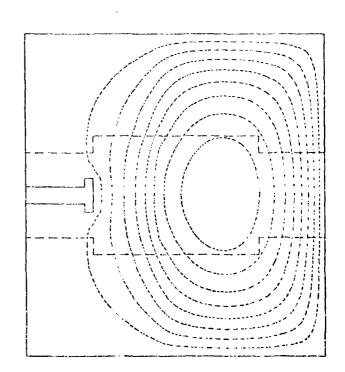
PLOT 20



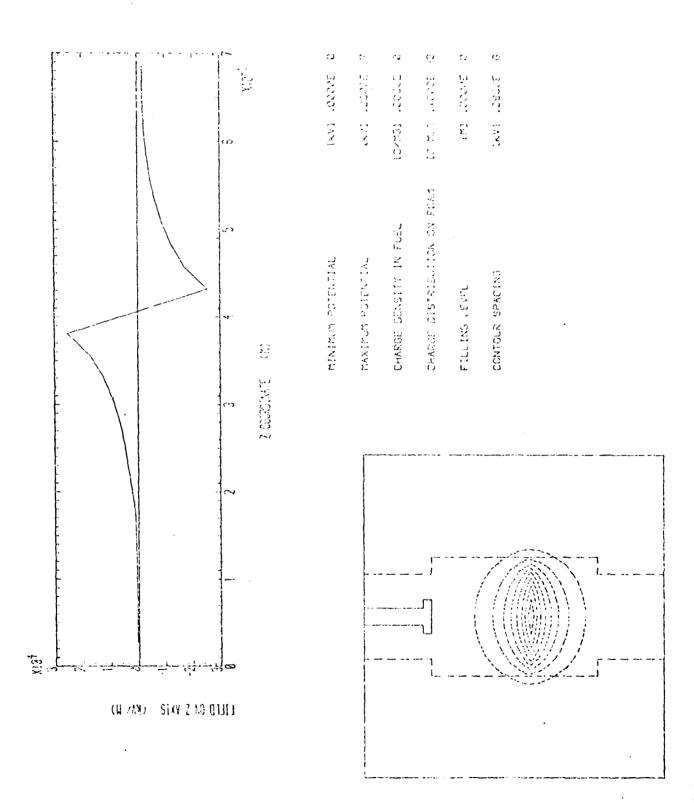
PLOT 21



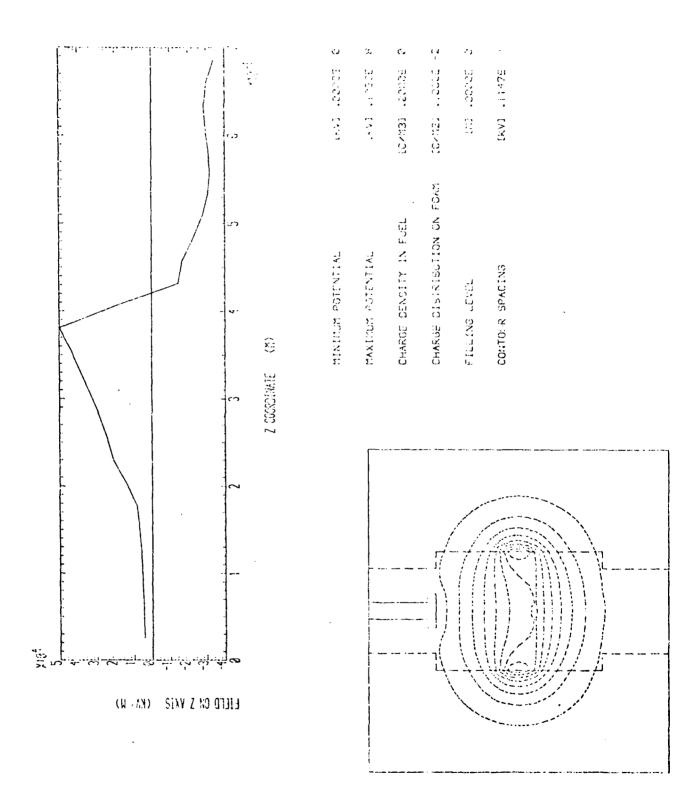




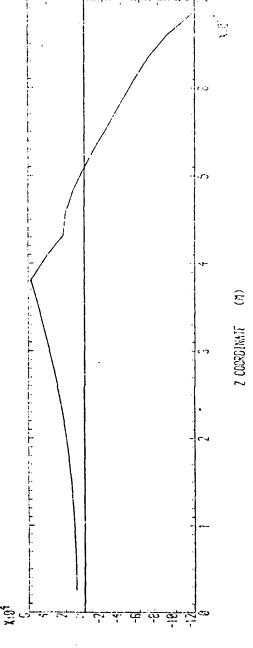
PLOT 22

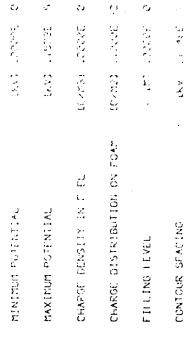


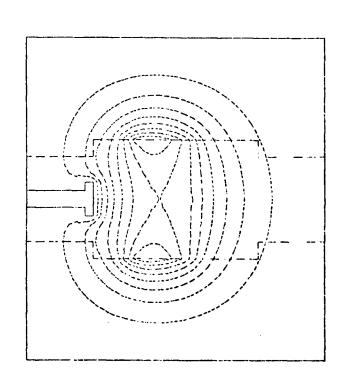
PLOT 23



PLOT 24

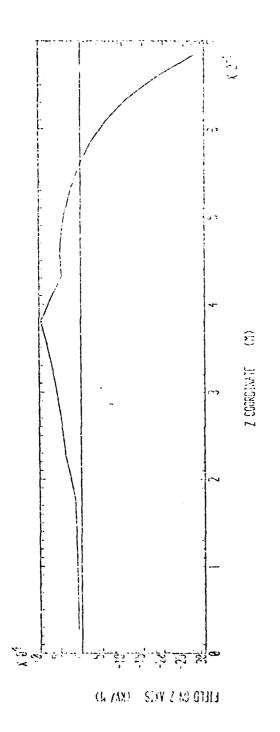


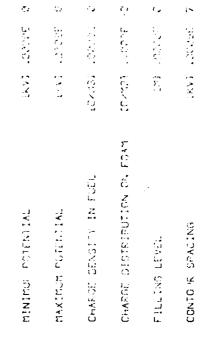


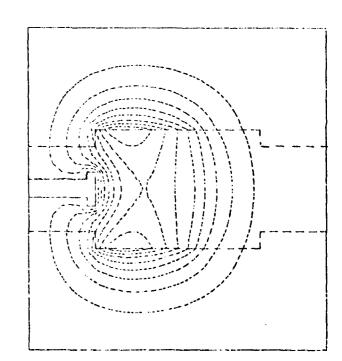


CHILL ON Z AXIS (RAVID)

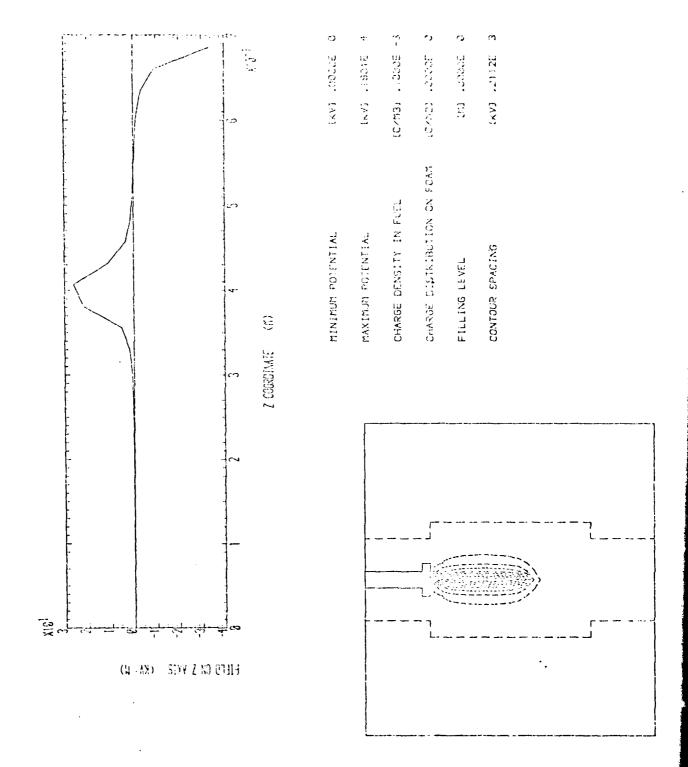
PLOT 25



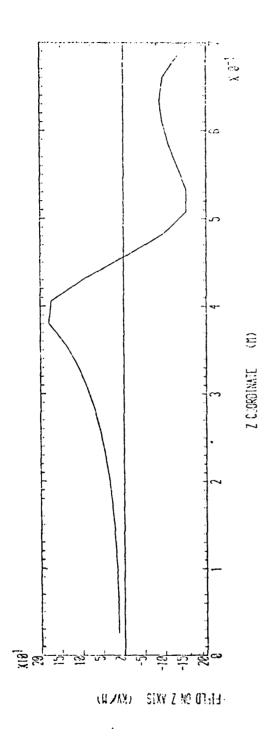


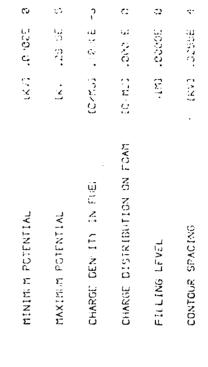


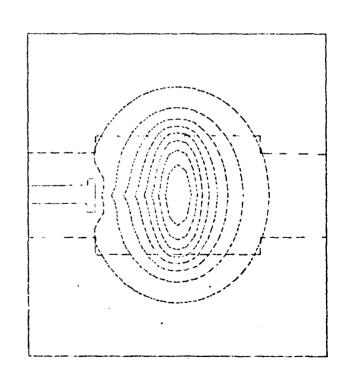
PLOT 26



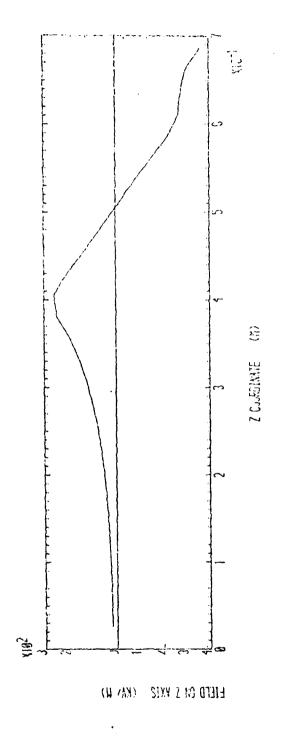
PLOT 27

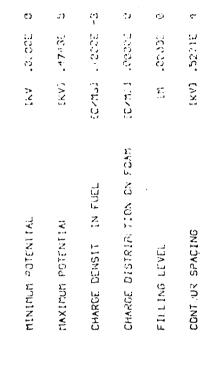


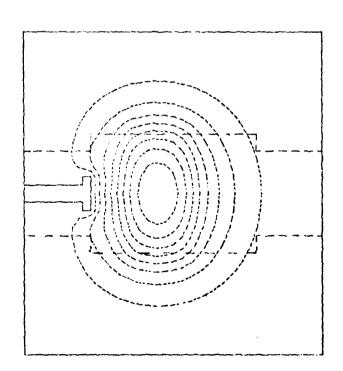




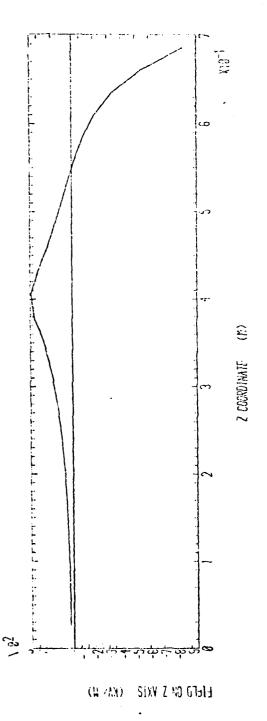
PLOT 28

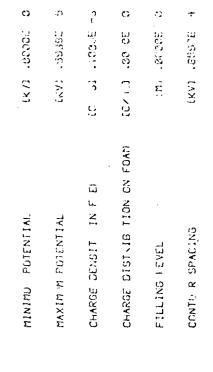


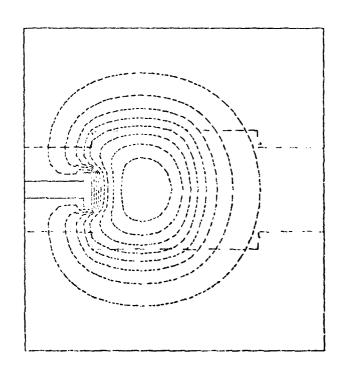




PLOT 29







PLOT 30

PLOTS 31 - 39

The configuration for plots 31-39 consists of foam Section 4 together with foam Sections 1, 2, and 3 inserted. (See Figure 8).

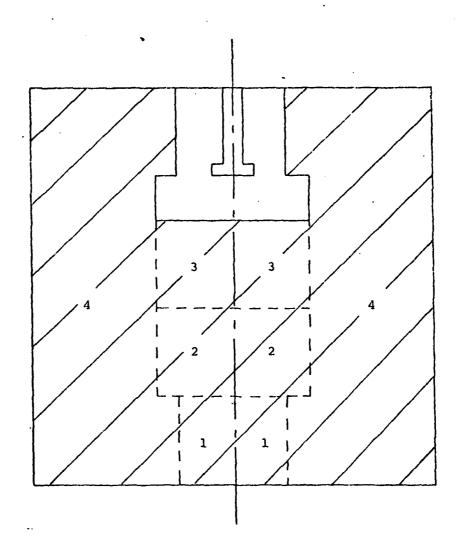
Plots 31 - 35	Charge	density	in	fuel			=	10-4	C/m ³		
	Charge	density	on	foam	sui	rface	#	c	C/m ²		
	Filling	levels	at	.1,	.2,	.3,	.4,	.5 n	etres	above	base.

Plot 36	Surface	change	density	of 10 ⁻³	C/m ²	on	Section	3
	upper su	ırface d	only.					

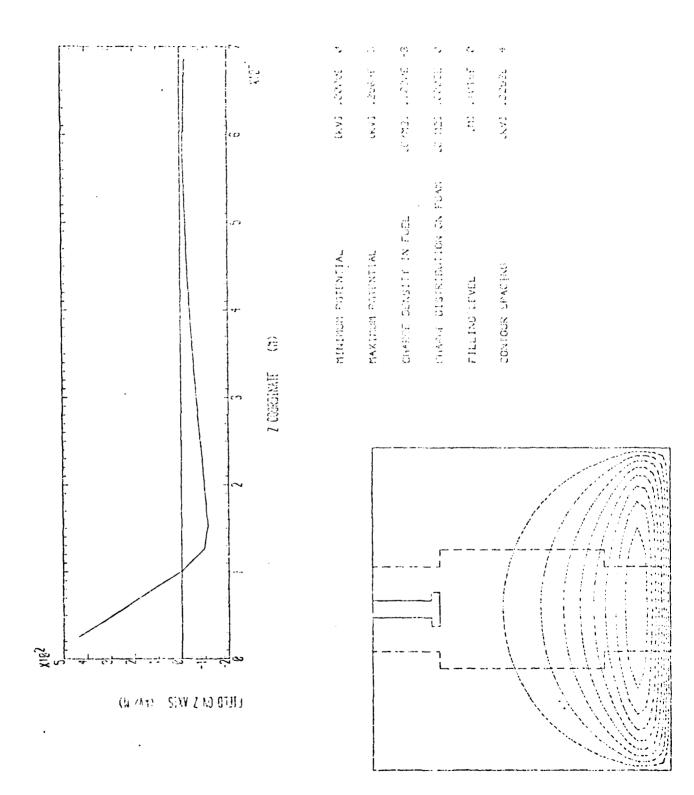
Plot 37	Surface	charge	on	Section	on 3	upper	surface	and	at
	height	.1 metre	e al	ove th	nis s	surface	≥.		

Plot 38	Vertical stream of fuel with charge density 10 4	C\w ₃
	impinging on the target area.	

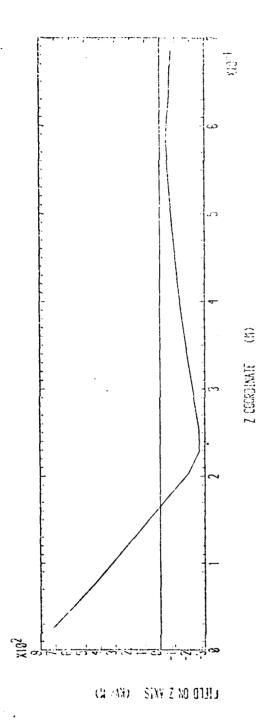
Plot 39 Voiding region filled with fuel.

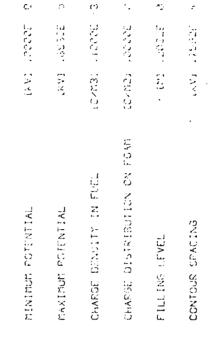


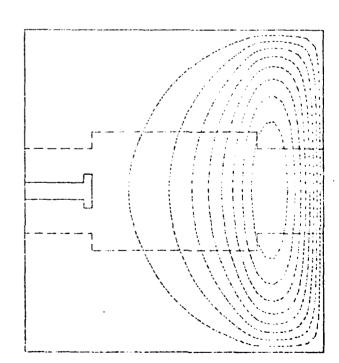
Foam sections 1 , 2 , 3 , and 4 in place fig. 8.



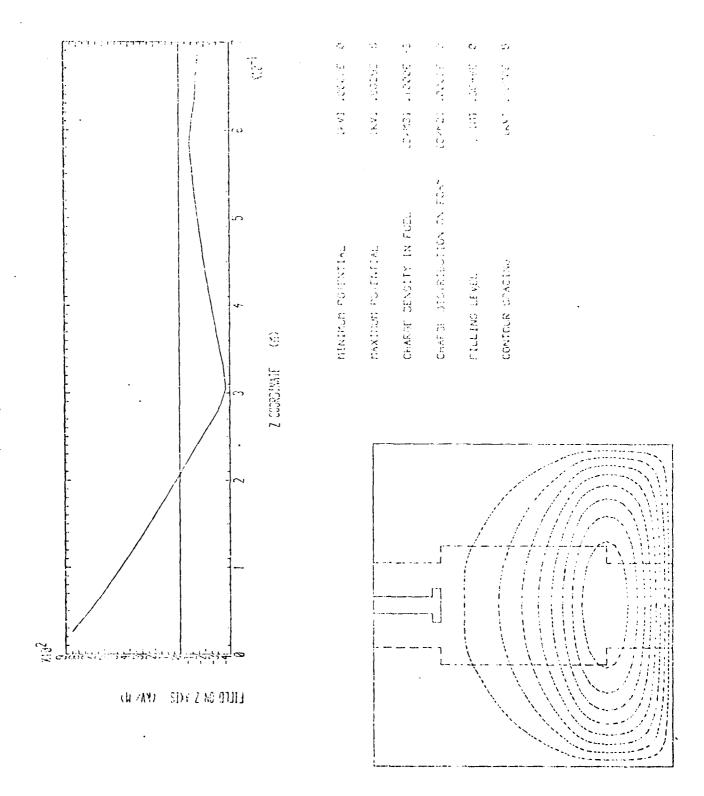
PLOT 31



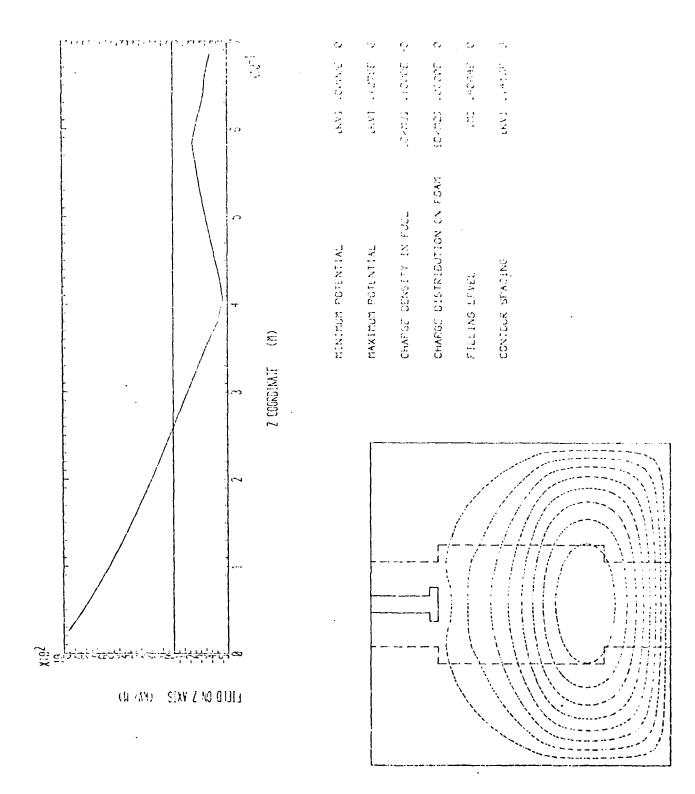




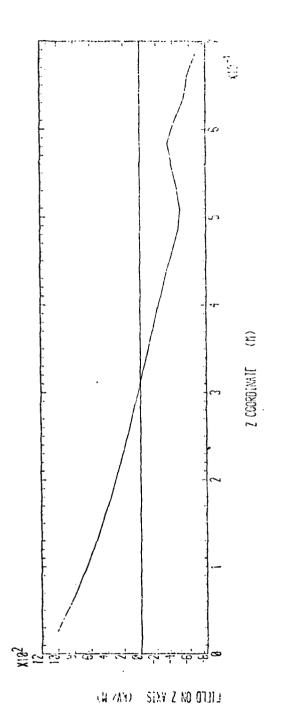
PLOT 32

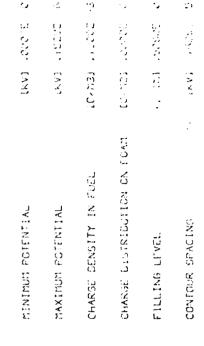


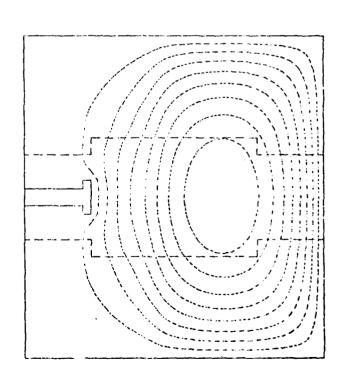
PLOT 33



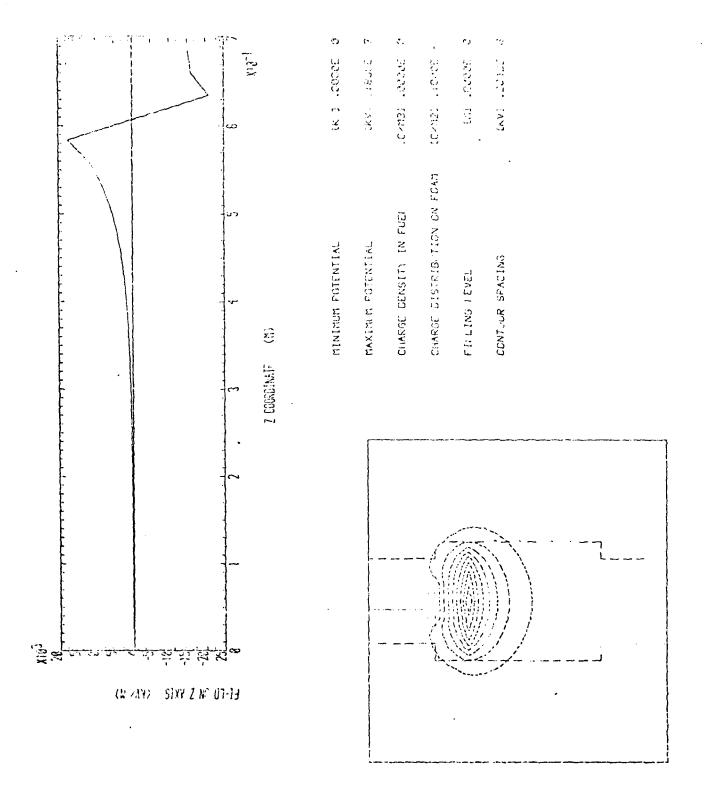
PLOT -34

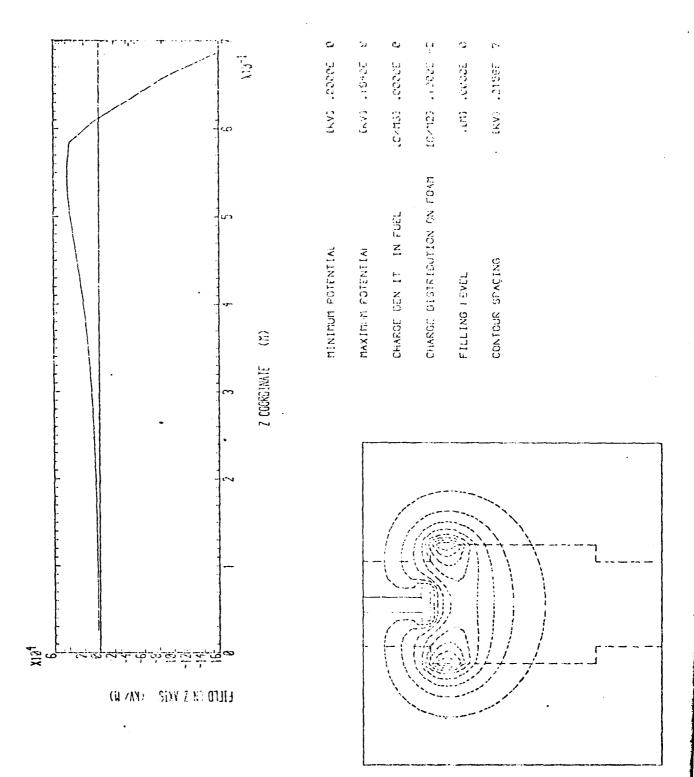




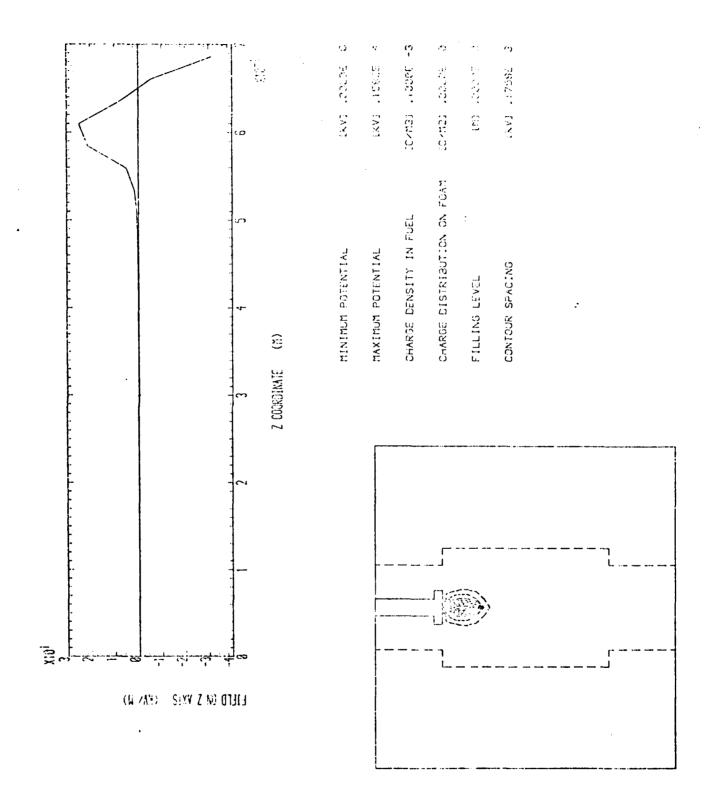


PLOT 35

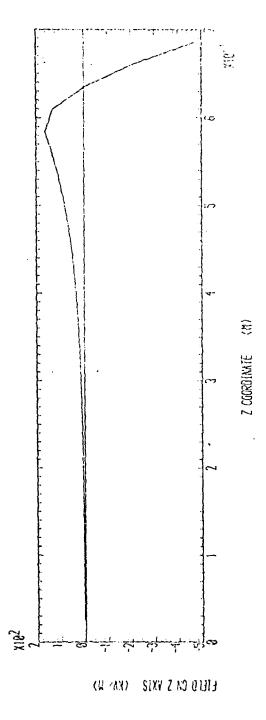


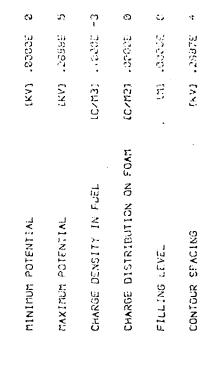


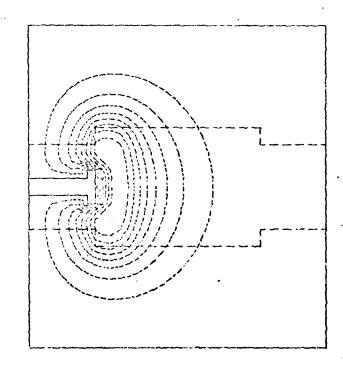
PLOT 37



PLOT 38





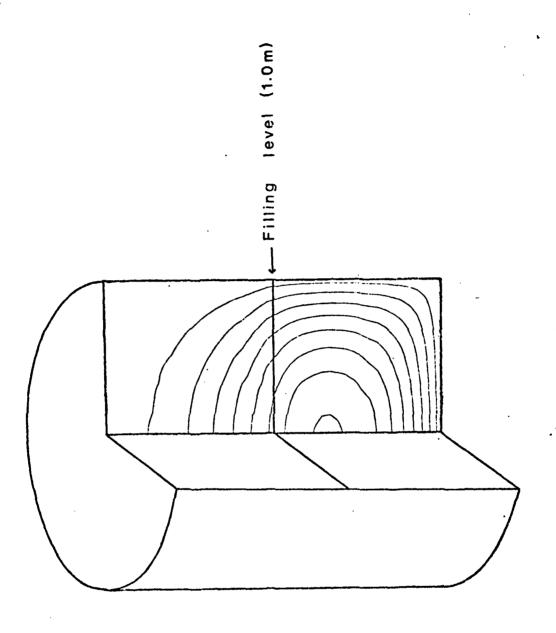


PLOT 39

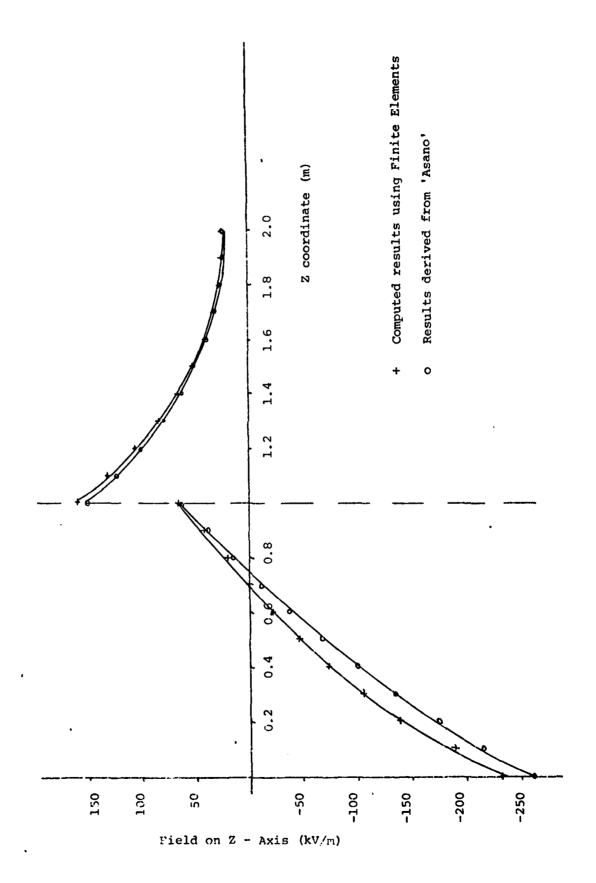
APPENDIX A

PROGRAM VALIDATION

The finite element program used in this study was written in FORTRAN. To test the accuracy of the program, it was checked against a problem for which an analytical solution has been found. Asano (Ref. 6) has provided an analytical solution for the potential in a cylindrical earthed metal tank half filled with charged liquid, assuming a uniform charge density. This problem was solved using the finite element program with a mesh of 760 nodes, and an equipotential map was drawn (Figure A]). The parameter chosen for comparison was the field along the cylinder axis. As can be seen from Figure A2, the finite element program produced results in good agreement with those of Asano.



Metal tank
Radius - 1.0 m
Height - 2.0 m



Validation using 'Asano' results

Fig. A2

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- E. Radgowski and R. Albrecht, "Investigation of electrostatic discharge in aircraft fuel tanks during refuelling", J. Aircraft, Vol. 16, No. 7, 1979 pp. 506-512.
- 2. S.J. Vellenga, "Estimating the electric field inside a rectangular tank with boundaries at zero potential", Appl. Sci. Res. 9, Section B.
- 3. J.A. Carruthers and K.J. Wigley, "The estimation of electrostatic potentials, fields and energies in a rectangular metal tank containing charged fuel", J. Inst. Petroleum, 48, 1962 pp. 181-189.
- 4. Electrostatic Test Summary, Aero Propulsion Laboratory, Wright-Patterson Air Force Base, 1980.
- 5. J.T. Leonard and W.A. Affens, "Electrostatic Charging of JP-4 fuel on polyurethane foams", NRL Report 8204, March 1978.
- K. Asano, "Electrostatic potential and field in a cylindrical tank containing charged liquid", Proc. IEE, Vol. 124, No. 12, 1977 pp. 1277-1281.